

Inhalt: 1. Besprechung des 4. Blattes

2. Komplexe Ginzburg-Landau-Gleichung

↳ Benjamin-Feir-Instabilität

3. Fitzhugh-Nagumo-Modell

4. Schlögl-Modell

1. Besprechung des 4. Blattes

↳ Lösung in Ausschnitt

↳ Code Beispiel

2. Komplexe Ginzburg-Landau Gleichung

$$ID: \frac{\partial}{\partial t} W(x,t) = W(x,t) - (1 + ic_2) |W(x,t)|^2 W(x,t) + (1 + ic_1) \frac{\partial^2}{\partial x^2} W(x,t)$$

uniforme/konvergente Oszillationen: $Q=0$

aus allgemeiner Ansatz für ebene Wellen: $W(x,t) = \alpha_Q \exp[i(\omega_Q t + \alpha x)]$

mit $|\alpha_Q|^2 = 1 - Q^2$ und $\omega_Q = -c_2 + (c_2 - c_1) Q^2$ folgt:

$$|\alpha_0|^2 = 1, \omega_0 = -c_2$$

$$\Rightarrow \text{Störung: } W \approx [1 + u(x,t)] e^{i\omega_0 t}$$

Ersatz in linearisierten Näherung:

$$\frac{\partial}{\partial t} u = -(1 + ic_2)(u + u^*) + (1 + ic_1) \frac{\partial^2}{\partial x^2} u$$

Weiter mit Fourier-Transformation:

$$u(x,t) = \int_{-\infty}^{\infty} dq u_q(t) e^{iqx}, \quad u^*(x,t) = \int_{-\infty}^{\infty} dq u_q^*(t) e^{-iqx}$$

$$\Rightarrow \frac{\partial u_q}{\partial t} = -(1 + ic_2)(u_q + u_q^*) - (1 + ic_1) q^2 u_q$$

$$\text{und } \frac{\partial u_q^*}{\partial t} = -(1 - ic_2)(u_q + u_q^*) - (1 - ic_1) q^2 u_q^*$$

2. Fitzhugh-Nagumo Modell

$$\varepsilon \dot{u} = u - \frac{u^3}{3} - v$$

$$\dot{v} = u + a$$

Aktivator (Schnell)

Inhibitor (Langsam)

a : Bifurkationsparameter

ε : Zeitskalentrennung $\varepsilon \ll 1$

Siehe Blatt 5

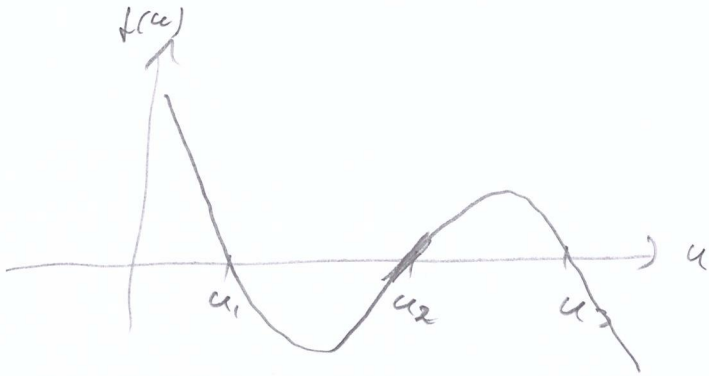
History of FHN

from Papers!

4. Schlögl-Modell:

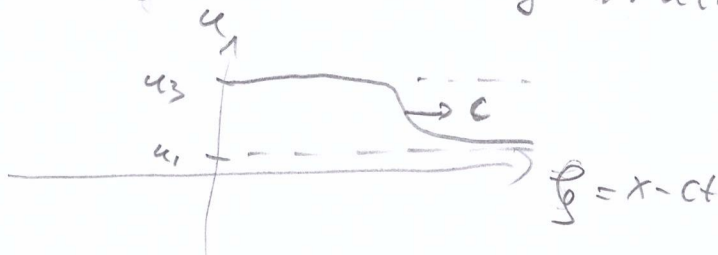
Dynamik von $u(x,t)$ gegeben durch $\frac{\partial}{\partial t} u(x,t) = f(u) + D \frac{\partial^2}{\partial x^2} u(x,t)$

mit $f(u) = k(u-u_1)(u-u_2)(u-u_3)$



Grechewitz-Ansatz

Mit bewegtes Koordinatensystem: $u(x,t) = u(\underbrace{x-ct}_{=\xi}) = u(\xi)$



$$\frac{\partial}{\partial t} \mapsto \frac{\partial}{\partial \xi} = \frac{\partial}{\partial (-ct)} \Rightarrow \frac{\partial}{\partial t} \mapsto -c \frac{\partial}{\partial \xi}, \quad \frac{\partial}{\partial x} \mapsto \frac{\partial}{\partial \xi}$$

$$\frac{\partial \xi}{\partial t} = -c$$

$$\Rightarrow \frac{\partial}{\partial t} u(x,t) = f(u) + D \frac{\partial^2}{\partial x^2} u(x,t)$$

$$-c u'(\xi) = f(u(\xi)) + D \frac{\partial^2}{\partial \xi^2} u(\xi)$$

$$u(\xi \rightarrow \pm \infty) = \begin{cases} u_1 \\ u_3 \end{cases} \quad \text{und} \quad \left. \frac{du}{d\xi} \right|_{\pm \infty} = 0$$

$$\Rightarrow \frac{du}{d\xi} = \beta (u-u_1)(u-u_3) - \frac{D}{c} \frac{d^2 u}{d\xi^2}$$

$$\frac{d\beta}{ds^2} = \frac{d}{ds} (\beta(u-u_1)(u-u_3)) = \beta \frac{d}{ds} (u^2 - u(u_1+u_3) + u_1u_3)$$

$$= \beta [2u u' - u'(u_1+u_3)]$$

$$= 2\beta u (\beta(u-u_1)(u-u_3)) - \beta (u-u_1)(u-u_3)(u_1+u_3)$$

$$\Rightarrow = (u-u_1)(u-u_3) [2\beta^2 u - \beta^2 (u_1+u_3)]$$

$$\rightarrow 0 = C u' + D u^2 + f(u)$$

$$= (u-u_1)(u-u_3) [C\beta + D 2\beta^2 u - D\beta^2 (u_1+u_3) - k(u-u_2)]$$

$$= (u-u_1)(u-u_3) [C\beta - \beta^2 D (u_1+u_3) - k u_2 + u(D 2\beta^2 - k)]$$

$u = u_1$ und $u = u_3$ Nullstellen!

3. Nullstelle: $[...] = 0$

$$\Rightarrow \beta = \sqrt{\frac{k}{2D}} \Rightarrow C \sqrt{\frac{k}{2D}} - \frac{k}{2D} D (u_1+u_3) - k u_2 = 0$$

$$\Rightarrow C = \sqrt{\frac{Dk}{2}} (u_1+u_3 - 2u_2)$$

$$u(\beta) = \frac{1}{2} [u_1+u_3 + (u_1-u_3) \operatorname{tanh} \left(\frac{1}{2} \sqrt{\frac{k}{2D}} (u_3-u_1) \xi \right)]$$