

6.2 Properties of inertia tensor (continued)

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$$(a) \quad \Theta_{jk} = \sum_{i=1}^N m_i \left[d^{(i)2} \delta_{jk} - d_j^{(i)} d_k^{(i)} \right]$$

$$(b) \quad \Theta_{jk} = \int d^3r \, \rho(\mathbf{r}) \left(d^2 \delta_{jk} - d_j d_k \right)$$

• diagonalization of $\underline{\Theta}$: $\underline{x} \rightarrow \underline{x}' = \underline{R} \underline{x}$

$$\underline{\Theta} \rightarrow \underline{\Theta}' = \underline{R}^T \underline{\Theta} \underline{R}$$

seek that $\underline{\Theta}' = \begin{pmatrix} \Theta_1 & 0 & 0 \\ 0 & \Theta_2 & 0 \\ 0 & 0 & \Theta_3 \end{pmatrix}$ Θ_i : principal moments of inertia
 $\underline{\omega}^{(i)}$: principal axes

$\Rightarrow$ eigenvalue equation: $\underline{\Theta} \underline{\omega}^{(i)} = \Theta_i \underline{\omega}^{(i)}$

$\Theta_1 = \Theta_2 = \Theta_3$ spherical top (symmetric, not necessarily a sphere!)

$\Theta_1 = \Theta_2 \neq \Theta_3$ symmetric top

$\Theta_1 \neq \Theta_2 \neq \Theta_3$ asymmetric top

• moment of inertia w.r.t. $\underline{\omega} = \underline{\hat{n}} \omega$: $\Theta_n = \underline{\hat{n}}^T \underline{\Theta} \underline{\hat{n}}$

$$= \sum_{i=1}^N m_i \left[d^{(i)2} - (\underline{d}^{(i)} \cdot \underline{\hat{n}})^2 \right]$$

$$= \sum_{i=1}^N m_i \left(\ell^{(i)} \right)^2$$

$$\ell^{(i)} = \left| \underline{d}^{(i)} - \underline{\hat{n}} (\underline{d}^{(i)} \cdot \underline{\hat{n}}) \right|$$

$\Rightarrow T_{\text{rot}} = \frac{1}{2} \Theta_n \omega^2$

• Steiner's theorem: $\Theta'_{jk} = \Theta_{jk} + M (a^2 \delta_{jk} - a_j a_k)$

with rotation axes of S and S' shifted by $\underline{a}$

• $\Theta_n' = \Theta_n + M \ell^2$, $\ell^2 = a^2 - (\underline{a} \cdot \underline{\hat{n}})^2$, $\Theta_i = \underline{\hat{n}}^T \underline{\Theta} \underline{\hat{n}}$

6.3 Angular momentum & equations of motion

Feb

$$(a) \underline{L} = \underbrace{M \underline{r}_0 \times \dot{\underline{r}}_0}_{\text{center of mass}} + \underbrace{\sum_{i=1}^N m_i \underline{d}^{(i)} \times (\underline{\omega} \times \underline{d}^{(i)})}_{\text{relative angular momentum}}$$

$\underline{L}_s$: center of mass $\underline{L}_R$: relative angular momentum

$$(b) \underline{L} = M \underline{r}_0 \times \dot{\underline{r}}_0 + \int d^3v \, \rho(\underline{r}) \, \underline{d} \times (\underline{\omega} \times \underline{d})$$

$$\Rightarrow \underline{L}_R = \underline{\Theta} \underline{\omega}$$

in principle: $\underline{L}_R \neq \underline{\omega}$ (rotation by $\underline{\Theta}$)

Wähle raumfestes System im Schwerpunkt: $\underline{r}_0 = 0$

$$\Rightarrow \underline{L} = \sum_{i=1}^N m_i \underline{d}^{(i)} \times (\underline{\omega} \times \underline{d}^{(i)}) = \underline{L}_R$$

$$\left(\text{Schwerpunkt: } \underline{r}_0 = \frac{1}{M} \sum_{i=1}^N m_i \underline{d}^{(i)} \right) \Rightarrow M \underline{r}_0 = \sum_{i=1}^N m_i \underline{d}^{(i)}$$

Newton'sche Gleichung für i -tes Teilchen: (vgl. 2.1)

$$m_i \underline{\ddot{d}}^{(i)} = \underbrace{\underline{F}^{(i)}_{\text{außen}}}_{\text{äußere Kräfte z.B. } m_i g} + \underbrace{\sum_{j=1}^N \underline{F}^{(ij)}}_{\text{innere Kräfte}}$$

$$\Rightarrow \underbrace{\sum_{i=1}^N m_i \underline{\ddot{d}}^{(i)}}_{M \underline{\ddot{r}}_0} = \sum_{i=1}^N \underline{F}^{(i)}_{\text{außen}} + \sum_{i,j=1}^N \underline{F}^{(ij)}$$

mit actio = reactio ($\underline{F}^{(ij)} = -\underline{F}^{(ji)} \Rightarrow \underline{F}^{(ij)} + \underline{F}^{(ji)} = 0$) folgt:

Schwerpunktssatz: $M \underline{\ddot{r}}_0 = \sum_{i=1}^N \underline{F}^{(i)}_{\text{außen}} =: \underline{F}^{(\text{außen})}$

Der Schwerpunkt bewegt sich so, als ob die gesamte Masse als starrer Körper in ihm versammelt wäre und die **Resultierende** $\underline{F}^{(\text{außen})}$ aller äußeren Kräfte in ihm angreift.

• innere Kräfte beeinflussen nicht den Schwerpunktsbewegung!

$$\underline{L} = \sum_{i=1}^N m_i \underline{d}^{(i)} \times \underline{\ddot{d}}^{(i)} = \sum_{i=1}^N \underline{d}^{(i)} \times \underline{F}^{(i)}_{\text{außen}} + \underbrace{\sum_{i,j=1}^N \underline{d}^{(i)} \times \underline{F}^{(ij)}}_{=0}$$

$$\frac{d}{dt} \underline{L}$$

$$\underline{F}^{(ij)} = -\underline{F}^{(ji)}$$

$$= 0$$

$$\underline{d}^{(i)} \times \underline{F}^{(ij)} + \underline{d}^{(j)} \times \underline{F}^{(ji)} = (\underline{d}^{(i)} - \underline{d}^{(j)}) \times \underline{F}^{(ij)}$$

Entlang Verbindungsvektor
zwischen i und $j \rightarrow \underline{d}^{(i)} - \underline{d}^{(j)}$

$$\Rightarrow \text{Drehimpulsatz: } \sum_{i=1}^N m_i \underline{\dot{d}}^{(i)} \times \underline{\ddot{d}}^{(i)} = \sum_{i=1}^N \underline{\dot{d}}^{(i)} \times \underline{\dot{d}}^{(i)} \text{ verschwindet}$$

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$$\Leftrightarrow \frac{d}{dt} \underline{L} = \underline{M} \quad \text{Drehmoment}$$

$$(i) \frac{d}{dt} \underline{L} = \frac{d}{dt} \left(\sum_{i=1}^N m_i \underline{\dot{d}}^{(i)} \times (\underline{\omega} \times \underline{\dot{d}}^{(i)}) \right) = \frac{d}{dt} (\underline{L}_i \underline{e}_i) = \frac{d}{dt} \sum_{i=1}^N \underline{\Theta}_{ij} \omega_j \underline{e}_i$$

$$= \sum_{j=1}^N \underline{\dot{\Theta}}_{ij} \omega_j \underline{e}_i + \sum_{j=1}^N \underline{\Theta}_{ij} \dot{\omega}_j \underline{e}_i$$

$$= \underbrace{\sum_{j=1}^N \underline{\dot{\Theta}}_{ij} \omega_j \underline{e}_i}_{\text{Abh. f. Beobachter}} + \sum_{j=1}^N \underline{\Theta}_{ij} \omega_j \underbrace{(\underline{\omega} \times \underline{e}_i)}_{\substack{\text{Körperfeste} \\ \text{bleibt}}} = \frac{d}{dt} \underline{L} + \underline{\omega} \times \underline{L} = \underline{M}$$

Abh. f. Beobachter
des Körpersystems, $\underline{e}_i$ konstant

Körperfeste
bleibt

(ii) alternativer Treuesatz von Leibniz auf Körpersystemen ($\underline{L}_0 = 0$):

$$\frac{d}{dt} = \left(\frac{d}{dt} \right)^* + \underline{\omega} \times$$

$$\Rightarrow \frac{d}{dt} \underline{L} = \left(\frac{d}{dt} \right)^* \underline{L} + \underline{\omega} \times \underline{L} = \underline{M}$$

$$= \underline{\underline{\Theta}} \underline{\dot{\omega}} + \underline{\omega} \times \underline{\underline{\Theta}} \underline{\omega} = \underline{M}$$

$\uparrow$
fest im Körpersystem

$\Rightarrow$ Euler'sche Gleichungen: Körperfeste Achsen = Hauptträgheitsachsen: $\underline{L}_i = \underline{\Theta}_{ii} \omega_i$

$$\Theta_1 \dot{\omega}_1 + (\Theta_3 - \Theta_2) \omega_2 \omega_3 = M_1$$

$$\Theta_2 \dot{\omega}_2 + (\Theta_1 - \Theta_3) \omega_1 \omega_3 = M_2$$

$$\Theta_3 \dot{\omega}_3 + (\Theta_2 - \Theta_1) \omega_1 \omega_2 = M_3$$

$$\Rightarrow \underline{\omega} \times \underline{\Theta} \underline{\omega} = \begin{pmatrix} \omega_1 \\ \omega_2 \\ \omega_3 \end{pmatrix} \times \begin{pmatrix} \Theta_1 \omega_1 \\ \Theta_2 \omega_2 \\ \Theta_3 \omega_3 \end{pmatrix}$$

• nichtlinear, gekoppelt!

6.4. Beispiele
 Bsp. 1) Symmetrischer Kegel: $\Theta_1 = \Theta_2 = \Theta \neq \Theta_3$ (kugelförmig) 77
 im Körperfesten System S^* (Struktur anfolgenden weggelassen)

$$\Rightarrow \Theta_3 \dot{\omega}_3 \neq (\Theta - \Theta) \omega_2 \omega_3 = 0 \Rightarrow \omega_3 = \text{const.}$$

$$\Theta_1 \dot{\omega}_1 = (\Theta - \Theta_3) \omega_2 \omega_3 \Rightarrow \dot{\omega}_1 = \frac{\Theta - \Theta_3}{\Theta} \omega_2 \omega_3$$

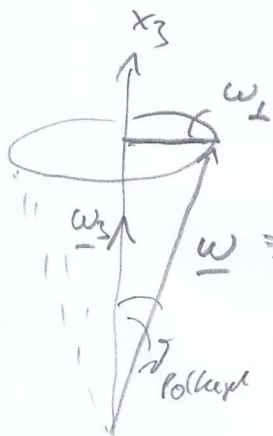
$$\Rightarrow \ddot{\omega}_1 = \frac{\Theta - \Theta_3}{\Theta} \omega_3 \dot{\omega}_2 = - \left(\frac{\Theta - \Theta_3}{\Theta} \omega_3^2 \right) \omega_1$$

$$\dot{\omega}_2 = - \frac{\Theta - \Theta_3}{\Theta} \omega_1 \omega_3$$

$$\Rightarrow \omega_1(t) = \omega_{\perp} \sin(\omega_0 t - \varphi_0) \text{ mit } \omega_0 = \frac{\Theta - \Theta_3}{\Theta} \omega_3$$

$$\omega_2(t) = -\omega_{\perp} \cos(\omega_0 t - \varphi_0), \quad \omega_1, \varphi_0 \text{ Integrationskonstanten}$$

$$\Rightarrow \omega_1^2 + \omega_2^2 = \omega_{\perp}^2 = \text{const.}$$



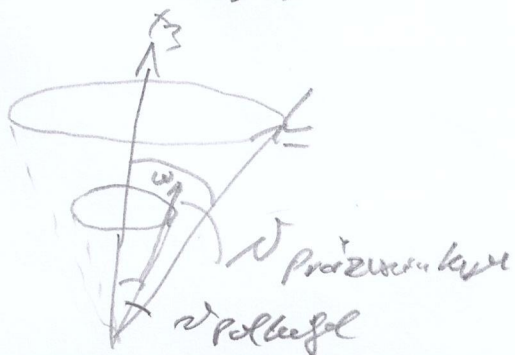
$$\Rightarrow \omega_1^2 + \omega_2^2 + \omega_3^2 = \omega_{\perp}^2 + \omega_3^2 \Rightarrow \vartheta_{\text{Polkegel}} = \arctan \frac{\omega_{\perp}}{\omega_3}$$

$\underline{\omega}$ und $\underline{L}$ (mit $L_j = \Theta_j \omega_j$) rotieren um
 Fluggeschwindigkeit

$$L^2 = \Theta^2 \omega_{\perp}^2 + \Theta_3^2 \omega_3^2 \quad (= L_{\perp}^2 + L_3^2)$$

$$\vartheta_{\text{Präzessionskegel}} = \arctan \frac{L_{\perp}}{L_3} = \arctan \frac{\Theta \omega_{\perp}}{\Theta_3 \omega_3} = \vartheta \text{ (Euler-Winkel)}$$

i.d. $\vartheta_{\text{Präzessionskegel}} \neq \vartheta_{\text{Polkegel}}$



Bislang: Körperfestes System! S^*

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Nun: Raumfestes System: S

Euler-Winkel:
$$\begin{pmatrix} \omega_1^* \\ \omega_2^* \\ \omega_3^* \end{pmatrix} = \dot{\varphi} \begin{pmatrix} \sin\varphi \sin\psi \\ \cos\varphi \sin\psi \\ \cos\psi \end{pmatrix} + \dot{\psi} \begin{pmatrix} \cos\varphi \\ -\sin\varphi \\ 0 \end{pmatrix} + \dot{\chi} \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix}$$

$$\begin{bmatrix} \omega_1 \\ \omega_2 \\ \omega_3 \end{bmatrix} = \dot{\varphi} \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix} + \dot{\psi} \begin{pmatrix} \cos\varphi \\ \sin\varphi \\ 0 \end{pmatrix} + \dot{\chi} \begin{pmatrix} \sin\varphi \sin\psi \\ -\cos\varphi \sin\psi \\ \cos\psi \end{pmatrix}$$

Nun: $\dot{\chi} = 0 \Rightarrow$
$$\begin{pmatrix} \omega_1^* \\ \omega_2^* \\ \omega_3^* \end{pmatrix} = \dot{\varphi} \begin{pmatrix} \sin\varphi \sin\psi \\ \cos\varphi \sin\psi \\ \cos\psi \end{pmatrix} + \dot{\psi} \begin{pmatrix} \cos\varphi \\ \sin\varphi \\ 0 \end{pmatrix}$$

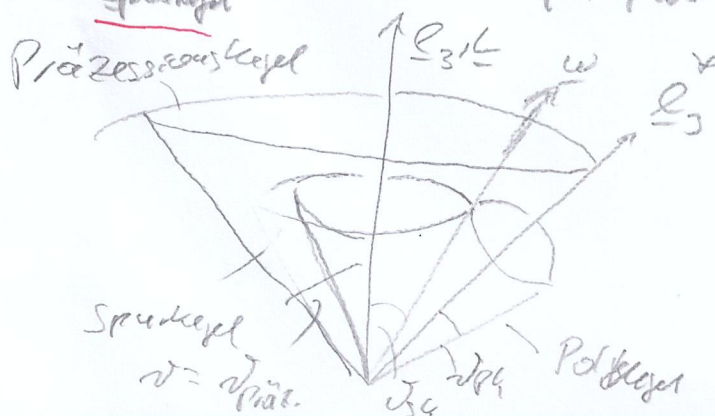
$$= \begin{pmatrix} \omega_3^* \sin(\psi(t)) \\ \omega_3^* \cos(\psi(t)) \\ \cos\psi \end{pmatrix}$$

$$\Rightarrow \varphi = \Omega t \Rightarrow \dot{\varphi} = \Omega = \text{const}$$

$$\dot{\varphi} = \frac{\omega_1^*}{\sin\psi} = \text{const} = \frac{\partial_s \omega_3^*}{\partial_{\cos\psi}} = \Omega' = \frac{L}{\Theta}$$

$\Rightarrow$ Euler-Winkel: $\psi(t) = \psi_0$, $\varphi(t) = \Omega t + \varphi_0$, $\chi(t) = \Omega' t + \chi_0$

Spinnkegel: $\frac{\dot{\varphi} \sin\psi}{\dot{\varphi} + \dot{\psi} \cos\psi} = \text{const} = \frac{\Omega' \sin\psi}{\Omega' + \Omega \cos\psi}$



wann der Erde
Kräftefreiheit sgam. Kugel:

$$\frac{\Omega - \Omega_3}{\Omega} \approx \frac{1}{300}$$

in Realität: $\frac{1}{430}$ (Bezeichnen...)